

Structure-guided 3D anisotropic tomography: a case study

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Abstract

Macro-velocity model building remains a challenge for seismic imaging. Currently, seismic datasets with large offsets and multi-azimuth usually require anisotropic models to produce a consistent 3D image of subsurface. The larger number of parameters to represent anisotropic models increases the ambiguity in 3D tomography. We present a successful case study illustrating the importance of using carefully designed conditioning of common image point gathers (CIP), dense RMO events picking and interactive QC tools, combined with 3D anisotropic structure guided tomography in order to reduce ambiguity and estimate a geologically conformable velocity model.

Introduction

Imaging seismic data with large offsets and multi-azimuths accurately demands 3D anisotropic tomography. The methodology for tomography in anisotropic media is well established in the literature [\(Barbosa et al.,](#page-3-0) [2008;](#page-3-0) [Bakulin](#page-3-1) [et al.,](#page-3-1) [2010;](#page-3-1) [Zhou et al.,](#page-4-0) [2011;](#page-4-0) [Wang and Tsvankin,](#page-4-1) [2013\)](#page-4-1), however, its implementation is not trivial. A carefully designed 3D tomography workflow still is an important and robust processing tool for 3D seismic imaging. We present a case study where two implementations of tomography produced noticeable differences in the velocity model and the resulting 3D seismic image. The baseline velocity model for our study is VTI and was produced by a service company after five iterations of 3D anisotropic tomography. The interpreter was not satisfied with the results and suggested further investigation for reducing the RMO through isotropic iterations. In order to evaluate the interpreter conjecture we used an in house implementation of 3D anisotropic structure-oriented tomography. The main features of our tomography implementation consist of a raytracing implementation for arbitrary Hamiltonians, model representation using nonuniform B-splines of variable order, structure-guided preconditioning [\(Hale,](#page-3-2) [2009\)](#page-3-2) and multi-scale iterations strategy. Also critically important was a robust implementation of dense event picking and interactive QC. Using this approach we validated the interpreter conjecture with a velocity model that reduced the RMO in all CIP compared with the baseline result for the whole offset range. Consequently, the resulted image is better focused. Surprisingly, the isotropic velocity

model, structurally conformable, was more consistent with the dataset than the baseline anisotropic model.

Methodology

The anisotropic 3D tomography follows closely the methodology for raytracing and computation of Frechét derivatives presented by [Barbosa et al.](#page-3-0) [\(2008\)](#page-3-0). Given an Hamiltonian $\mathcal{H} (\mathbf{x}, \mathbf{p}; \mathbf{C}) = 0$, defined by an eikonal equation for each position x and slowness vector p , for a set of material parameters represented by C. For example, to model qP TTI anisotropy we use the anelliptical approximation [\(Schoenberg and de Hoop,](#page-4-2) [2000\)](#page-4-2):

$$
\mathcal{H}(\mathbf{x}, \mathbf{p}) = \frac{1}{2} \left[C_{\varepsilon}(\mathbf{x}) (\mathbf{p} \cdot \mathbf{p} - (\mathbf{p} \cdot \mathbf{v}(\mathbf{x}))^2) + C_0(\mathbf{x}) (\mathbf{p} \cdot \mathbf{v}(\mathbf{x}))^2 - C_0(\mathbf{x}) (C_{\varepsilon}(\mathbf{x}) - C_{\delta}(\mathbf{x}) (\mathbf{p} \cdot \mathbf{p} - (\mathbf{p} \cdot \mathbf{v}(\mathbf{x}))^2) (\mathbf{p} \cdot \mathbf{v}(\mathbf{x}))^2 - 1 \right] = 0,
$$
\n(1)

where C_0 represents the square of propagation velocity along the symmetry axis, $C_{\varepsilon} \equiv C_0(1+2\varepsilon)$ and $C_{\delta} \equiv C_0(1+2\varepsilon)$ 2δ); ε and δ are the Thomsen's parameters [\(Thomsen,](#page-4-3) [1986\)](#page-4-3); $v(x)$ indicates the direction of the symmetry axis.

The residual moveout at CIP gathers relative to a reference offset h_0 , i.e., $\delta z_{\text{rmo}} \equiv z(h) - z(h_0)$, can be modeled by the fundamental equation of tomography in the migrated domain:

$$
z(h) - z(h_0) = -\frac{1}{\|\mathbf{p}^s(h) + \mathbf{p}^r(h)\|\mathbf{n} \cdot \mathbf{e}_z}
$$

$$
\left(\int_0^{T^s(h)} \frac{\partial \mathcal{H}}{\partial \mathbf{C}} \delta \mathbf{C} d\tau + \int_0^{T^r(h)} \frac{\partial \mathcal{H}}{\partial \mathbf{C}} \delta \mathbf{C} d\tau\right)
$$

$$
+ \frac{1}{\|\mathbf{p}^s(h_0) + \mathbf{p}^r(h_0)\|\mathbf{n} \cdot \mathbf{e}_z}
$$

$$
\left(\int_0^{T^s(h_0)} \frac{\partial \mathcal{H}}{\partial \mathbf{C}} \delta \mathbf{C} d\tau + \int_0^{T^r(h_0)} \frac{\partial \mathcal{H}}{\partial \mathbf{C}} \delta \mathbf{C} d\tau\right). (2)
$$

Equation [\(2\)](#page-0-0) determines a linear relationship between the event residual moveout and the perturbations of velocity model parameters δC around a current reference model C. In order to build a linear system using this equation we need to compute pairs of specular rays from each picked event to the surface constrained to fit the offsets, h and h_0 , and azimuth; p^s and p^r are the initial slowness vector at CIP depth for the ray branch to source and to receiver, respectively; *T ^s* and *T r* represent traveltimes for each branch from the event coordinates to the surface for a defined offset and azimuth, n is the direction normal to a possible reflector and, e*z* the vertical direction.

The model parameters are specified using B-splines interpolation [\(De Boor et al.,](#page-3-3) [1978;](#page-3-3) [Piegl and Tiller,](#page-3-4) [2012\)](#page-3-4). The value of each parameter at a given position, $C^n(x, y, z)$ for $n \in \{1,...,N\}$, is uniquely defined by corresponding

 coefficients C_{ijk}^n specified in a 3D control mesh of dimensions $N_x \times N_y \times N_z$, i.e.,

$$
C^{n}(x,y,z) = \sum_{i}^{N_{x}} \sum_{j}^{N_{y}} \sum_{k}^{N_{z}} C_{ijk}^{n} B_{1}^{i}(x) B_{2}^{j}(y) B_{3}^{k}(z) ,
$$
 (3)

where $B_1^i(x), B_2^j(y), B_3^k(z)$ represents the B-splines polynomial bases along each coordinate direction. Substituting the representation above in equation [\(2\)](#page-0-0) produces the linear system which allows us to determine δC_{ijk}^n and update the model. The implementations permits nonuniform meshes and B-Splines of variable order.

The linear system resulting from equation [\(2\)](#page-0-0) is, in general, poorly conditioned [\(Woodward et al.,](#page-4-4) [2008;](#page-4-4) [Golub and](#page-3-5) [Van Loan,](#page-3-5) [2012\)](#page-3-5). In order to assure stable linear iterations it is necessary the introduction of supplementary information about the model properties. A classical approach to stabilization is regularization [\(Costa et al.,](#page-3-6) [2008\)](#page-3-6). More recently, in order to estimate geologically conformable models, inclusion of structural information to better conditioning geophysical inversion have been and continue being developed [\(Clapp et al.,](#page-3-7) [2004;](#page-3-7) [Ma et al.,](#page-3-8) [2012;](#page-3-8) [Shank et al.,](#page-4-5) [2017;](#page-4-5) [Li et al.,](#page-3-9) [2018;](#page-3-9) [Zhou et al.,](#page-4-6) [2016;](#page-4-6) [Guo et al.,](#page-3-10) [2017\)](#page-3-10). Our implementation used the structural smoothing equation [\(Hale,](#page-3-2) [2009\)](#page-3-2) to enforce a conformable geological model,

$$
m(\mathbf{x}) - \frac{1}{2} \nabla \cdot (\alpha(\mathbf{x}) \mathbf{D}(\mathbf{x}) \nabla m(\mathbf{x})) = q(\mathbf{x}), \qquad (4)
$$

where $m(x)$ represents the update of a model component, $\delta C^n({\bf x})$, and $q({\bf x})$ its reparameterization following [Harlan](#page-3-11) [\(1995\)](#page-3-11); $D(x)$ is the diffusivity tensor, i.e., the inverse of the structure tensor [\(Fehmers and Höcker,](#page-3-12) [2003\)](#page-3-12), finally $\alpha(x)$ is a normalized coherence and allows for structural discontinuities[\(Hale,](#page-3-2) [2009\)](#page-3-2). The numerical solution of equation [\(4\)](#page-1-0) in a dense grid formally corresponds to a model reparameterization, which can be represented by the matrix operation $m = Pq$. The model update requires the least squares solution of the linear system:

$$
\left(\begin{array}{c}\mathbf{LP} \\ \lambda_m\mathbf{I}\end{array}\right)\mathbf{q}=\left(\begin{array}{c}\mathbf{d} \\ \mathbf{0}\end{array}\right)\,,\tag{5}
$$

where L represents the tomographic matrix associated to equation [\(2\)](#page-0-0), d contains selected RMOs and λ*m* is a damping parameter. We use the LSMR algorithm [\(Fong](#page-3-13) [and Saunders,](#page-3-13) [2011\)](#page-3-13) to solve the linear system in equation [\(5\)](#page-1-1).

Figura 1: Dense picking of RMO events and interactive QC.

RMO picking and quality control are critical for successful field data applications of seismic tomography. The event selection is performed in three steps: 1) free picking of the largest number of events up to a prescribed number; 2) selection of events with automatic quality control and, 3) fitting of the selected events with a polynomial of variable degree. Figure [1](#page-1-2) shows and example of a CIP gather with the curves resulting from the steps in the event selection, the red curves represent the first step in the selection process, orange curves are produced by the automatic quality control, cyan curves correspond to the polynomial fitting of selected events. One can notice that the polynomial fitting also contributes to QC discarding some of the previously selected events.

Figure contains the flowchat of our methodology for structure-guided anisotropic tomography.

Figura 2: Flowchat for structure-guided anisotropic 3D tomography.

Field data applications

Initially we present field data tomography results to stress the importance of the model preconditioning in our methodology. Figures [3](#page-2-0)[-6](#page-2-1) present, in a single color scale, the model updates using four types preconditioning by model reparameterization overlaid by the corresponding migrated image. All the remaining parameters indicated in the flowchart in Figure [2](#page-1-3) remain the same. We claim that the fault in the middle of this section present higher definition in Figure 6 , also the velocity model update is more conformable to the seismic image. This feature of structure-guided tomography can be very important for seismic interpretation.

8 iterações LSMR Figura 3: Model update without preconditioning overlaid to the corresponding migrated image.

8 iterações LSMR

Figura 4: Model update without preconditioning and using exponential filter overlaid with the corresponding migrated image.

8 iterações LSMR

Figura 5: Model update with preconditioning by exponential smoothing overlaid with the corresponding migrated image.

8 iterações LSMR

Figura 6: Model update with preconditioning by structural smoothing overlaid with the corresponding migrated image.

Our next application of the structure-guided tomography to a field dateset presents a case study where two implementations of tomography produced noticeable differences in the estimated velocity models and the resulting 3D seismic image. The baseline result is an VTI velocity model produced by a service company after five iterations of a 3D anisotropic tomography. Figure [8](#page-3-14) shows a selected section extracted from the 3D model estimated by VTI tomography overlaid with the corresponding migrated image. The interpreter was not satisfied with the results and suggested further investigation for reducing the RMO through isotropic iterations. In order to evaluate the interpreter conjecture we used the workflow for 3D anisotropic structure-oriented tomography, as presented in the methodology section. We performed two iterations of structure-guided 3D isotropic tomography. The initial model was the vertical velocity produced by the VTI tomography. For the first iteration we used a more rigorous quality control and constrained the maximum velocity update to 25m/s and a grid of 150 $m \times 150$ m \times 25 m. For the second iteration the quality control was relaxed and the maximum velocity update was set to 300m/s and used refined grid 50 $m \times 50$ m \times 10 m. Figure [7](#page-3-15) present a CIP gather after each iteration. One can notice the improvement in the flatness of the CIP produced by the model update in our second iteration over the whole depth range, which corroborates the interpreter hypothesis for this area.

Figures and show, respectively, the stacked migrated section for the original VTI model and the isotropic model after two iterations of structure-guided tomography overlaid with the corresponding velocity model. The isotropic model improved image focusing all over this section. The isotropic velocity model, structurally conformable, was more consistent with this dataset than the initial VTI model.

Figura 7: CIP gathers after two tomography iterations.

Figura 8: Vertical velocity model derived from VTI tomography overlaid with the migrated image.

Discussion

Our implementation of structure-guided tomography rely on dense event picking and interactive QC, a robust strategy to extract structural dip information and conditioning of the diffusivity tensor field and, preconditioning using the smoothing equation proposed by [Hale](#page-3-2) [\(2009\)](#page-3-2). Though the preconditioning at each iteration of model update requires the numerical solution of a differential equation, our numerical experiments show that it reduces significantly the number of iteration required for the convergence of the linear system solver. Moreover, this preconditioning strategy produced better results enforcing structural information on the velocity updates than alternative, computationally less expensive, filtering and smoothing operators. Another feature that was important for applications to field data sets is the model representation, using B-splines of variable order and nonuniform meshes allows the implementation of multi-scale iterations.

We presented a successful application of structureguided 3D tomography to improve seismic imaging and contribute to the interpretation in a field dataset offshore southeast Brazil. This case study illustrates how structureguided isotropic iterations can fit large offset data with a structure conformable velocity model. The multi-scale strategy was instrumental to improve the resolution of the model updates. The final velocity model produced flatter CIP gathers at all depth and offset range and, significantly improved the focusing of reflection events in the final stacked migrated image. Last but not least, this result shows once again that facing ambiguity in seismic interpretation requires, besides better tools to enforce geological constraints in velocity model building, a close collaboration with the interpretation group.

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